

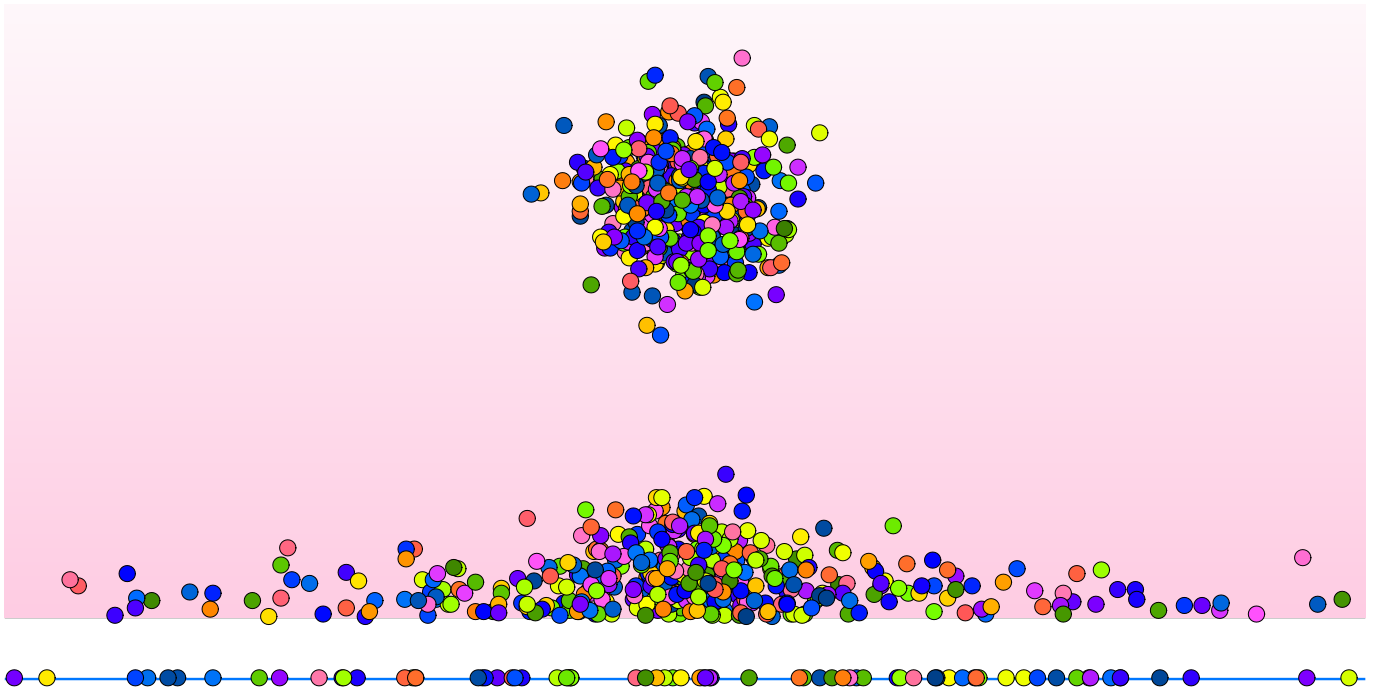
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# MAT<sub>CH</sub>A

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(MAThematical models for fast diffusion CHAnnels in population dynamics and epidemiology)

## Annexe II



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## Fujita on a Heat-exchanger-system

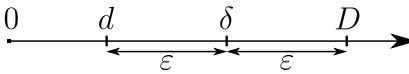
### I.1 The linear model

We consider in this section the following *Heat-exchanger-system*

$$\begin{cases} \partial_t u = d\Delta u - u + v, & t > 0, \quad x \in \mathbb{R} \\ \partial_t v = D\Delta v + u - v, & t > 0, \quad x \in \mathbb{R}. \end{cases} \quad (\text{I.1})$$

Without loss of generality, one assumes that  $D \geq d$  and one sets for convenience:

$$\delta := \frac{D+d}{2}, \quad \text{and} \quad \varepsilon := \frac{D-d}{2}.$$

Notice  $d > 0$  requires  $\delta > \varepsilon$ . 

**Remark.** If  $\varepsilon = 0$ , we have  $d = D = \delta$ ; whence, calling  $\sigma := u + v$  and  $w := u - v$ , one gets the uncoupled system

$$\begin{cases} \partial_t \sigma = \delta \Delta \sigma, & t > 0, \quad x \in \mathbb{R} \\ \partial_t w = \delta \Delta w - 2w & t > 0, \quad x \in \mathbb{R} \end{cases}$$

whose straight resolution yields

$$u(t, x) = \frac{1 + e^{-2t}}{2} [G_\delta(t, \bullet) * u_0](x) + \frac{1 - e^{-2t}}{2} [G_\delta(t, \bullet) * v_0](x)$$

$$v(t, x) = \frac{1 - e^{-2t}}{2} [G_\delta(t, \bullet) * u_0](x) + \frac{1 + e^{-2t}}{2} [G_\delta(t, \bullet) * v_0](x).$$

It is then clear that  $\|u(t, \bullet)\|_\infty, \|v(t, \bullet)\|_\infty \lesssim t^{-1/2}$ .

In a more general way the latter move does not work; however, we are able to provide an explicit expression of the Fourier transform of the solution:

$$\begin{aligned}\widehat{u}(t, \xi) &= \frac{(\sqrt{s^1} + \varepsilon\xi^2)e^{-t(1-\sqrt{s^1})} + (\sqrt{s^1} - \varepsilon\xi^2)e^{-t(1+\sqrt{s^1})}}{2\sqrt{s^1}} e^{-\delta t\xi^2} \widehat{u}_0(\xi) \\ &\quad + \frac{e^{-t(1-\sqrt{s^1})} - e^{-t(1+\sqrt{s^1})}}{2\sqrt{s^1}} e^{-\delta t\xi^2} \widehat{v}_0(\xi) \\ \widehat{v}(t, \xi) &= \frac{e^{-t(1-\sqrt{s^1})} - e^{-t(1+\sqrt{s^1})}}{2\sqrt{s^1}} e^{-\delta t\xi^2} \widehat{u}_0(\xi) \\ &\quad + \frac{(\sqrt{s^1} - \varepsilon\xi^2)e^{-t(1-\sqrt{s^1})} + (\sqrt{s^1} + \varepsilon\xi^2)e^{-t(1+\sqrt{s^1})}}{2\sqrt{s^1}} e^{-\delta t\xi^2} \widehat{v}_0(\xi)\end{aligned}$$

with  $\sqrt{s^1} := \sqrt{1 + \varepsilon^2\xi^4}$ . As in the latter remark, it seems that  $u$  and  $v$  can be split into a *persistent part* and a *residual part*:

$$u(t, x) = u_\infty(t, x) + u_r(t, x) \quad \text{and} \quad v(t, x) = v_\infty(t, x) + v_r(t, x),$$

with

$$\begin{aligned}\widehat{u}_\infty(t, \xi) &= \frac{(\sqrt{s^1} + \varepsilon\xi^2)e^{-t(1-\sqrt{s^1})}}{2\sqrt{s^1}} e^{-\delta t\xi^2} \widehat{u}_0(\xi) + \frac{e^{-t(1-\sqrt{s^1})}}{2\sqrt{s^1}} e^{-\delta t\xi^2} \widehat{v}_0(\xi) \\ \widehat{u}_r(t, \xi) &= \frac{(\sqrt{s^1} - \varepsilon\xi^2)e^{-t(1+\sqrt{s^1})}}{2\sqrt{s^1}} e^{-\delta t\xi^2} \widehat{u}_0(\xi) - \frac{e^{-t(1+\sqrt{s^1})}}{2\sqrt{s^1}} e^{-\delta t\xi^2} \widehat{v}_0(\xi) \\ \widehat{v}_\infty(t, \xi) &= \frac{e^{-t(1-\sqrt{s^1})}}{2\sqrt{s^1}} e^{-\delta t\xi^2} \widehat{u}_0(\xi) + \frac{(\sqrt{s^1} - \varepsilon\xi^2)e^{-t(1-\sqrt{s^1})}}{2\sqrt{s^1}} e^{-\delta t\xi^2} \widehat{v}_0(\xi) \\ \widehat{v}_r(t, \xi) &= -\frac{e^{-t(1+\sqrt{s^1})}}{2\sqrt{s^1}} e^{-\delta t\xi^2} \widehat{u}_0(\xi) + \frac{(\sqrt{s^1} + \varepsilon\xi^2)e^{-t(1+\sqrt{s^1})}}{2\sqrt{s^1}} e^{-\delta t\xi^2} \widehat{v}_0(\xi).\end{aligned}$$

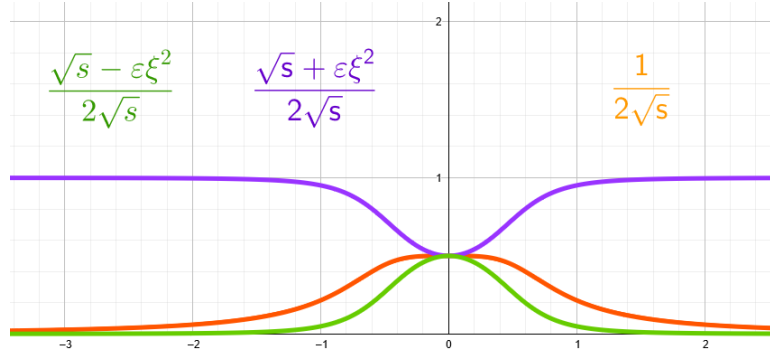
**Proposition 1 (The residual parts vanish exponentially fast)**

Assuming  $u_0$  and  $v_0$  to be smooth, there is  $c > 0$  such that  $\|u_r(t, \bullet)\|_\infty < ce^{-2t}$  and  $\|v_r(t, \bullet)\|_\infty < ce^{-2t}$ . Furthermore,  $c$  is explicitly known and only depends on  $N$  and the  $L^1$ -norms of  $\widehat{u}_0$  and  $\widehat{v}_0$ .

**Lemma 2 (Boundedness)**

One has,

$$\boxed{1} \quad 0 < \frac{\sqrt{s^1} - \varepsilon\xi^2}{2\sqrt{s^1}} \leq \frac{1}{2} \leq \frac{\sqrt{s^1} + \varepsilon\xi^2}{2\sqrt{s^1}} < 1, \quad \boxed{2} \quad 0 < \frac{1}{2\sqrt{s^1}} \leq \frac{1}{2},$$


**Proof (Lemma 2)**

We have  $\sqrt{s^1} = \sqrt{1 + \varepsilon^2\xi^4} \geq 1$ , hence [2] is clearly true; then for  $\xi \neq 0$ ,

$$0 < \frac{\sqrt{s^1} - \varepsilon\xi^2}{2\sqrt{s^1}} = \frac{1}{2} - \frac{1}{2\sqrt{1 + \frac{1}{\varepsilon^2\xi^4}}} < \frac{1}{2},$$

$$\frac{1}{2} < \frac{\sqrt{s^1} + \varepsilon\xi^2}{2\sqrt{s^1}} = \frac{1}{2} + \frac{1}{2\sqrt{1 + \frac{1}{\varepsilon^2\xi^4}}} < 1,$$

and for  $\xi = 0$ ,  $\frac{\sqrt{s^1} \pm \varepsilon\xi^2}{2\sqrt{s^1}} = \frac{1}{2}$ . That proves [1].  $\square$

**Proof (Proposition 1)**

It is sufficient to show the controls for the  $L^1$ -norm of  $\hat{u}(t, \bullet)$  and  $\hat{v}(t, \bullet)$ , since

$$|u(t, x)| = \frac{1}{2\pi} \left| \int_{\mathbb{R}} \hat{u}(t, \xi) e^{i\xi x} d\xi \right| \leq \frac{1}{2\pi} \int_{\mathbb{R}} |\hat{u}(t, \xi)| d\xi = \frac{1}{2\pi} \|\hat{u}(t, \bullet)\|_{L^1(\mathbb{R})}.$$

We have  $1 + \sqrt{s^1} \geq 2$ , hence, using Lemma 2,

$$\int_{\mathbb{R}} |\hat{u}_r(t, \xi)| d\xi \leq e^{-2t} \int_{\mathbb{R}} \frac{1}{2} |\hat{u}_0(\xi)| + \frac{1}{2} |\hat{v}_0(\xi)| d\xi$$

$$\int_{\mathbb{R}} |\hat{v}_r(t, \xi)| d\xi \leq e^{-2t} \int_{\mathbb{R}} \frac{1}{2} |\hat{u}_0(\xi)| + |\hat{v}_0(\xi)| d\xi,$$

whence

$$\|u_r(t, \bullet)\|_{\infty} \leq e^{-2t} \left( \frac{1}{4\pi} \|\hat{u}_0\|_{L^1(\mathbb{R})} + \frac{1}{4\pi} \|\hat{v}_0\|_{L^1(\mathbb{R})} \right)$$

$$\|v_r(t, \bullet)\|_{\infty} \leq e^{-2t} \left( \frac{1}{4\pi} \|\hat{u}_0\|_{L^1(\mathbb{R})} + \frac{1}{2\pi} \|\hat{v}_0\|_{L^1(\mathbb{R})} \right). \quad \square$$

We are now investigating what kind of equation solve the persistent parts  $u_\infty$  and  $v_\infty$ . On the Fourier-side, we own

$$\hat{u}_\infty(t, \xi) = \left[ \frac{\sqrt{s^1} + \varepsilon \xi^2}{2\sqrt{s^1}} \hat{u}_0(\xi) + \frac{1}{2\sqrt{s^1}} \hat{v}_0(\xi) \right] e^{t(\sqrt{s^1} - \delta \xi^2 - 1)} =: \hat{u}_{0,\infty}(\xi) e^{t(\sqrt{s^1} - \delta \xi^2 - 1)}$$

$$\hat{v}_\infty(t, \xi) = \left[ \frac{1}{2\sqrt{s^1}} \hat{u}_0(\xi) + \frac{\sqrt{s^1} - \varepsilon \xi^2}{2\sqrt{s^1}} \hat{v}_0(\xi) \right] e^{t(\sqrt{s^1} - \delta \xi^2 - 1)} =: \hat{v}_{0,\infty}(\xi) e^{t(\sqrt{s^1} - \delta \xi^2 - 1)}$$

hence  $\hat{u}_\infty(\cdot, \xi)$  and  $\hat{v}_\infty(\cdot, \xi)$  satisfy the Cauchy problems ( $\xi$  plays as a parameter)

$$\begin{cases} \partial_t \hat{u}_\infty = (\sqrt{s^1} - \delta \xi^2 - 1) \hat{u}_\infty \\ \hat{u}_\infty|_{t=0} = \hat{u}_{0,\infty} \end{cases} \quad \begin{cases} \partial_t \hat{v}_\infty = (\sqrt{s^1} - \delta \xi^2 - 1) \hat{v}_\infty \\ \hat{v}_\infty|_{t=0} = \hat{v}_{0,\infty} \end{cases}$$

As a consequence, by taking the inverse Fourier transform, we are led to these *non-local* and *uncoupled* diffusion equations:

$$\begin{cases} \partial_t u_\infty = \mathcal{F}_\xi^{-1} \left[ \sqrt{1 + \varepsilon^2 \xi^4} - \delta \xi^2 \right] * u_\infty - u_\infty, & t > 0, \quad x \in \mathbb{R} \\ u_\infty|_{t=0} = u_{0,\infty} & x \in \mathbb{R} \end{cases}$$

$$\begin{cases} \partial_t v_\infty = \mathcal{F}_\xi^{-1} \left[ \sqrt{1 + \varepsilon^2 \xi^4} - \delta \xi^2 \right] * v_\infty - v_\infty, & t > 0, \quad x \in \mathbb{R} \\ v_\infty|_{t=0} = v_{0,\infty} & x \in \mathbb{R} \end{cases}$$

We set:

$$J(x) := \mathcal{F}_\xi^{-1} \left[ \sqrt{1 + \varepsilon^2 \xi^4} - \delta \xi^2 \right] (x).$$

It is clear that  $J$  is a probability kernel since

$$\|J\|_{L^1(\mathbb{R})} = \left[ \sqrt{1 + \varepsilon^2 \xi^4} - \delta \xi^2 \right] \Big|_{\xi=0} = 1.$$

Now, near the origin,

$$\hat{J}(\xi) = 1 - \delta \xi^2 + \frac{\varepsilon^2}{2} \xi^4 + o(\xi^4),$$

and, as  $|\xi| \rightarrow \infty$ ,

$$\hat{J}(\xi) = -d \xi^2 + \frac{1}{2\varepsilon \xi^2} + o\left(\frac{1}{\xi^2}\right).$$

